

# Solid-State Mesostructured Perovskite $\text{CH}_3\text{NH}_3\text{PbI}_3$ Solar Cells: Charge Transport, Recombination, and Diffusion Length

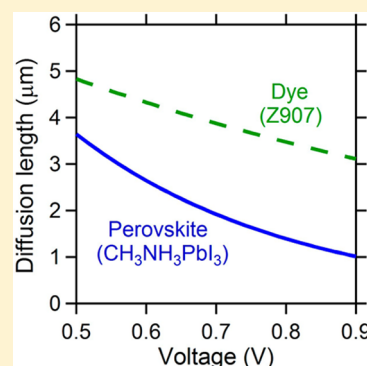
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**S** Supporting Information

**ABSTRACT:** We report on the effect of  $\text{TiO}_2$  film thickness on charge transport and recombination in solid-state mesostructured perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  (via one-step coating) solar cells using spiro-MeOTAD as the hole conductor. Intensity-modulated photocurrent/photovoltage spectroscopies show that the transport and recombination properties of solid-state mesostructured perovskite solar cells are similar to those of solid-state dye-sensitized solar cells. Charge transport in perovskite cells is dominated by electron conduction within the mesoporous  $\text{TiO}_2$  network rather than from the perovskite layer. Although no significant film-thickness dependence is found for transport and recombination, the efficiency of perovskite cells increases with  $\text{TiO}_2$  film thickness from 240 nm to about 650–850 nm owing primarily to the enhanced light harvesting. Further increasing film thickness reduces cell efficiency associated with decreased fill factor or photocurrent density. The electron diffusion length in mesostructured perovskite cells is longer than 1  $\mu\text{m}$  for over four orders of magnitude of light intensity.

**SECTION:** Energy Conversion and Storage; Energy and Charge Transport



Organometallic halide perovskites (e.g.,  $\text{CH}_3\text{NH}_3\text{PbI}_3$  and  $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Cl}_x$ ) have recently emerged as a new class of light absorbers that have demonstrated exceptional progress in solar cell performance.<sup>1,2</sup> Efficiencies over 15% have been reported<sup>3,4</sup> since their first use in solar cells in 2009.<sup>5</sup> Such high-efficiency perovskite cells can be made by either solution processing or thermal evaporation. These promising results have made halide perovskites one of the most attractive light-absorbing materials for solar conversion applications.<sup>6–16</sup> Two types of device architectures (i.e., mesostructured or bilayer thin-film solar cells) are normally used for halide perovskites. The first type is adapted from solid-state dye-sensitized solar cells (DSSCs), in which a mesoporous metal oxide (e.g.,  $\text{TiO}_2$ ) provides a large surface area to support the perovskite absorber and spiro-MeOTAD is normally used as the hole transport material (HTM).<sup>3,8,9,15,17,18</sup> The device also works well in the absence of the HTM layer.<sup>19</sup> In the second type, the bulk thin-film perovskite absorber layer is sandwiched between the electron- and hole-contact layers (e.g.,  $\text{TiO}_2$  and spiro-MeOTAD, respectively).<sup>4,20</sup> Despite the rapid progress in device performance, understanding the structural and electronic properties of halide perovskites is just in its infancy. One of the key determinants of solar cell performance is the charge-carrier diffusion length ( $L_d$ ). A longer  $L_d$  generally means a thicker absorber layer for greater light harvesting. Recent studies have shown that for the planar perovskite cell structure, both electron and hole diffusion lengths are  $>1 \mu\text{m}$  for  $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Cl}_x$ <sup>21</sup> and  $>100 \text{ nm}$  for  $\text{CH}_3\text{NH}_3\text{PbI}_3$ .<sup>22</sup> However, the charge-carrier diffusion length has not been determined for the solid-state mesostructured perovskite cells,

even though such a cell structure has led to the highest certified perovskite solar cell efficiencies so far. Because the solid-state mesostructured perovskite solar cell resembles the solid-state DSSC,<sup>23,24</sup> several studies report on the transport properties of mesostructured perovskite cells using common characterization techniques from the DSSC field, such as transient photocurrent/photovoltage decay<sup>9,15</sup> and impedance spectroscopy (IS).<sup>16,25,26</sup> One study shows that the transport time in mesostructured perovskite cells depends little on the hole mobility of HTM (varying over five orders of magnitude),<sup>15</sup> suggesting that transport in these cells is limited by electron conduction. Another study suggests that electron transport could take place in the mesoporous metal oxide layer (e.g.,  $\text{TiO}_2$ ) or in the perovskite layer itself.<sup>9</sup>

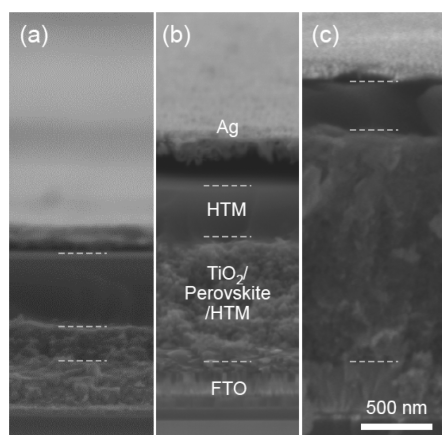
Herein we present results of our investigation on the effect of  $\text{TiO}_2$  film thickness on charge transport, recombination, and device characteristics of solid-state mesostructured perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solar cells using spiro-MeOTAD as the hole conductor. Charge transport and recombination are examined by intensity-modulated photocurrent/photovoltage spectroscopies (IMPS/IMVS). We find that the transport and recombination properties of solid-state mesostructured perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solar cells are similar to those of solid-state Z907-based DSSC. The electron diffusion length in mesostructured perovskite cells is longer than 1  $\mu\text{m}$  under normal cell operation conditions.

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Figure 1 shows typical cross-sectional scanning electron microscopy (SEM) images of solid-state mesostructured  $\text{TiO}_2/\text{CH}_3\text{NH}_3\text{PbI}_3/\text{spiro-MeOTAD}$  solar cells with different  $\text{TiO}_2$  film thicknesses. Three representative  $\text{TiO}_2$  film thicknesses are shown in panels a–c, respectively.



**Figure 1.** Typical cross-sectional SEM images of solid-state mesostructured  $\text{TiO}_2/\text{CH}_3\text{NH}_3\text{PbI}_3/\text{spiro-MeOTAD}$  solar cells with different  $\text{TiO}_2$  film thicknesses. Three representative  $\text{TiO}_2$  film thicknesses are shown in panels a–c, respectively.

$\text{CH}_3\text{NH}_3\text{PbI}_3/\text{spiro-MeOTAD}$  solar cells. The average  $\text{TiO}_2$  film thickness (as determined from a surface profiler) ranges from 240 to 1650 nm. For each sample in this study a spiro-MeOTAD overlayer is seen on the top surface of the mesoporous  $\text{TiO}_2$  layer. The thickness of this HTM capping layer decreases from about 500 to 350 nm when the  $\text{TiO}_2$  film thickness increases from 240 to 1650 nm. The thickness range of the HTM capping layer is consistent with literature reports.<sup>4,8,9,25</sup> The typical X-ray diffraction (XRD) patterns and ultraviolet–visible (UV–vis) absorption spectra (Figure S1 in the Supporting Information) for the perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  agree with our previous studies.<sup>27,28</sup> The absorption spectrum of the  $\text{CH}_3\text{NH}_3\text{PbI}_3$ -sensitized mesoporous  $\text{TiO}_2$  film increases for thicker  $\text{TiO}_2$  films, corresponding to an enhanced loading of  $\text{CH}_3\text{NH}_3\text{PbI}_3$ .

Table 1 compares the  $J$ – $V$  characteristics of solid-state mesostructured perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solar cells with

**Table 1. Effect of  $\text{TiO}_2$  Film Thickness on Short-Circuit Photocurrent Density ( $J_{\text{sc}}$ ), Open-Circuit Voltage ( $V_{\text{oc}}$ ), Fill Factor (FF), and Conversion Efficiency  $\eta$  of Solid-State Perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  Solar Cells**

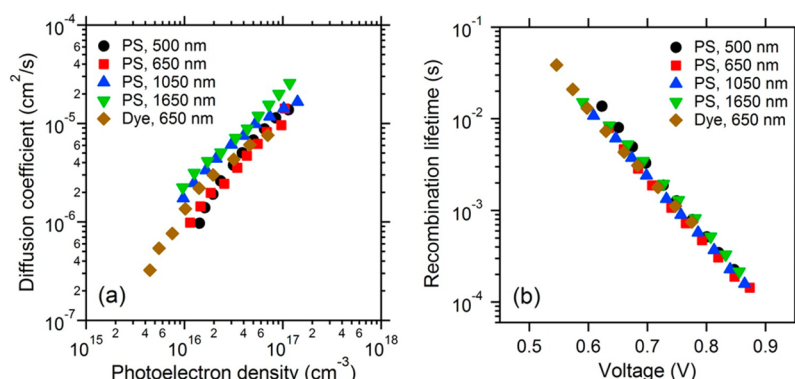
thickness (nm)	$J_{\text{sc}}$ (mA/cm <sup>2</sup> )	$V_{\text{oc}}$ (V)	FF	$\eta$ (%)
240	10.41	0.851	0.577	5.11
300	11.13	0.859	0.593	5.67
400	13.11	0.871	0.548	6.26
500	15.36	0.869	0.536	7.16
650	17.57	0.889	0.567	8.85
850	17.69	0.890	0.555	8.73
1050	17.16	0.878	0.490	7.39
1650	14.17	0.869	0.445	5.48

different  $\text{TiO}_2$  film thicknesses measured under simulated AM 1.5G 100 mW/cm<sup>2</sup>. The  $J$ – $V$  and IPCE curves as a function of  $\text{TiO}_2$  film thickness are, respectively, shown in Figures S2 and S3 (Supporting Information). The cell efficiency first increases from about 5 to 9% when the  $\text{TiO}_2$  film thickness is changed from 240 nm to the range of 650–850 nm. Further increasing film thickness results in lower cell efficiencies. The initial

increase in the cell efficiency with increasing film thickness is consistent with the enhanced light absorption (Figure S1 in the Supporting Information). The decrease in cell efficiency for  $\text{TiO}_2$  films thicker than 1  $\mu\text{m}$  is primarily caused by the reduced FF or  $J_{\text{sc}}$ , despite the enhanced absorption for thicker  $\text{TiO}_2$  films. The significant drop of FF for  $\text{TiO}_2$  films thicker than 1  $\mu\text{m}$  is likely caused by the reduced pore filling of spiro-MeOTAD within the mesoporous  $\text{TiO}_2$  films. It is worth noting that the maximum efficiency for a perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solar cell using a 650 nm  $\text{TiO}_2$  film in this study is 10.3% with a  $J_{\text{sc}}$  of 17.8 mA/cm<sup>2</sup>,  $V_{\text{oc}}$  of 0.90 V, and FF of 0.64 ( $J$ – $V$  and IPCE curves shown in Figure S4 in the Supporting Information). In contrast, a solid-state DSSC using the Z907 dye and 650 nm  $\text{TiO}_2$  film shows a  $J_{\text{sc}}$  of 4.5 mA/cm<sup>2</sup>,  $V_{\text{oc}}$  of 0.77 V, and FF of 0.63 to yield an efficiency of 2.2%. The efficiency difference between the perovskite cell and DSSC results primarily from their markedly different  $J_{\text{sc}}$  values associated with their respective light-absorption properties.<sup>6</sup>

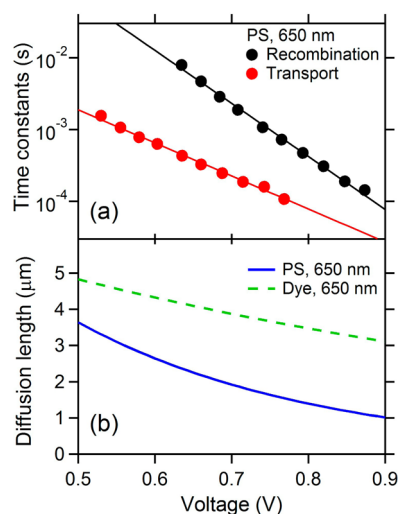
Figure 2a shows the effect of  $\text{TiO}_2$  film thickness on the electron diffusion coefficient ( $D$ ) for the solid-state mesostructured perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solar cells as a function of photoelectron density ( $n$ ). The  $D$  values are calculated from transport times ( $\tau_t$ ) determined by IMPS measurements at short circuit.<sup>29</sup> The diffusion coefficients for a Z907-sensitized solid-state dye cell based on the 650-nm  $\text{TiO}_2$  film are also plotted for comparison. The  $D$  values can be fitted to a power-law dependence on  $n$  following an expression  $D \propto n^{1/\alpha-1}$  with  $\alpha = 0.49$  for all samples. Such power-law dependence in DSSCs has been attributed to multiple trapping and detrapping of electrons during their transit through the  $\text{TiO}_2$  network.<sup>30–35</sup> There is no obvious thickness dependence of the  $D$  values among perovskite solar cells. The diffusion coefficients for the perovskite cell with a 650 nm  $\text{TiO}_2$  film are essentially the same as that for the solid-state DSSC with the same  $\text{TiO}_2$  film thickness. These results indicate that the mechanistic factors governing the charge transport in the solid-state  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solar cells are similar to those in the solid-state DSSCs, suggesting that transport in the perovskite cells is dominated by the electron conduction within the  $\text{TiO}_2$  network rather than from the perovskite layer itself. Figure 2b shows the recombination lifetime ( $\tau_r$ ) as a function of quasi-Fermi level (or open-circuit voltage). The  $\tau_r$  values are determined by IMVS measurements at open circuit.<sup>29</sup> All perovskite cells and the DSSC exhibit essentially the same voltage dependence of recombination lifetime. The  $\tau_r$  value decreases with increasing voltage for all samples; for DSSCs, such dependence is normally attributed to the enhanced recombination with the higher electron density.<sup>29</sup> For a given voltage, there is no noticeable difference of  $\tau_r$  among perovskite cells with different  $\text{TiO}_2$  film thickness, suggesting that recombination occurs homogeneously across the mesoporous  $\text{TiO}_2$  film.<sup>36,37</sup> Moreover, no difference in  $\tau_r$  is observed between the perovskite cells and DSSC, implying that recombination kinetics does not depend on the absorber deposited on the  $\text{TiO}_2$  surface in this study. This observation significantly differs from a recent study on solid-state perovskite cells using ZnO nanorods, where recombination is found to increase with the length of ZnO nanorods.<sup>38</sup>

From the measurements of transport and recombination times ( $\tau_t$  and  $\tau_r$ , respectively), one can determine the electron diffusion length  $L_d$  in the solid-state mesostructured perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solar cells by using the expression  $L_d = d(\tau_t/\tau_r)^{1/2}$  (where  $d$  is the  $\text{TiO}_2$  film thickness). The electron diffusion



**Figure 2.** (a) Electron diffusion coefficient as a function of photoelectron density. (b) Recombination lifetime as a function of open-circuit voltage. PS: perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$ ; Dye: Z907 dye.

length in DSSC is a measure of the average distance an electron travels before it recombines with either the oxidized absorber or the hole conductor.<sup>39</sup> A longer  $L_d$  could lead to higher charge-collection or light-harvesting efficiencies for an improved solar conversion efficiency. Figure 3a shows the dependence of



**Figure 3.** (a) Comparison of transport and recombination times for a mesostructured perovskite solar cell as a function of voltage (or quasi-Fermi level). (b) Comparison of the electron diffusion length for a mesostructured perovskite cell and a solid-state DSSC using 650 nm  $\text{TiO}_2$  films. PS: perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$ ; Dye: Z907 dye.

transport and recombination times on voltage (or quasi-Fermi level QFL) for a solid-state mesostructured perovskite solar cell using a 650 nm thick  $\text{TiO}_2$  film. The QFL at open circuit is given by the open-circuit voltage, whereas the average QFL at short circuit for a given light intensity is determined by a switching method reported previously for DSSCs.<sup>40</sup> Figure 3b shows the  $L_d$  values as a function of QFL for the perovskite solar cell. The diffusion length decreases from about 3.5 to 1  $\mu\text{m}$  when the QFL increases from 0.5 to 0.9 V (corresponding to over four orders of magnitude of light intensity). In comparison,  $L_d$  for a solid-state Z907-based DSSC exhibits a similar dependence on QFL, which is in agreement with a previous report on solid-state DSSCs.<sup>41</sup> Although  $L_d$  for the solid-state DSSC is longer than that for the perovskite cell over the entire voltage range, the lower absorption property of dye molecules compared with perovskites limits the performance of solid-state DSSCs. The observed voltage dependence of  $L_d$

implies that charge collection should, in general, be more efficient at lower light intensity for mesostructured perovskite cells. Because the average QFLs at short circuit (under simulated AM 1.5G 100  $\text{mW}/\text{cm}^2$ ) for perovskite cells are around 0.75 to 0.8 V in this study, the  $L_d$  values are thus expected to be about 1.2 to 1.5  $\mu\text{m}$  (Figure 3b), implying that charge collection at short circuit should not limit solar conversion efficiency for all perovskite samples except for the one using a 1.65  $\mu\text{m}$   $\text{TiO}_2$  film. This is consistent with the  $J_{sc}$  dependence on the  $\text{TiO}_2$  film thickness (Figure S2 in the Supporting Information). A poorer pore filling of spiro-MeOTAD for thicker  $\text{TiO}_2$  films could also contribute to the observed low performance (e.g., FF or  $J_{sc}$ ) for the perovskite cell based on thicker  $\text{TiO}_2$  films (Table 1).<sup>41</sup>

In summary, we examined charge transport, recombination, and device characteristics of solid-state mesostructured perovskite  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solar cells based on 0.24 to 1.65  $\mu\text{m}$  thick  $\text{TiO}_2$  films and spiro-MeOTAD hole conductor. Charge transport and recombination in the solid-state mesostructured perovskite cells are similar to those in the solid-state DSSC and exhibit little dependence on the  $\text{TiO}_2$  film thickness. The performance of perovskite cells increases with  $\text{TiO}_2$  film thickness up to 650–850 nm, resulting primarily from the enhanced light harvesting. Further increasing film thickness results in lower cell efficiencies, mainly caused by the reduced FF or  $J_{sc}$ . The electron diffusion length in mesostructured perovskite cells is found to be longer than 1  $\mu\text{m}$  under normal cell operation conditions.

## ■ ASSOCIATED CONTENT

### Supporting Information

Experimental method, XRD and UV-vis spectra,  $J$ – $V$  and IPCE curves as a function of  $\text{TiO}_2$  film thickness, optimized  $J$ – $V$  characteristics, and IPCE spectrum. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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### Notes

The authors declare no competing financial interest.

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