

Strain-symmetrized $\text{In}_x\text{Ga}_{1-x}\text{As}/\text{In}_y\text{Al}_{1-y}\text{As}$ HEMTs with extremely high 2DEG densities and mobilities

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ABSTRACT: Strain-symmetrized $\text{In}_x\text{Ga}_{1-x}\text{As}/\text{In}_y\text{Al}_{1-y}\text{As}$ hetero-structures with $0.53 \leq x \leq 0.74$ and $0.52 \leq y \leq 0.415$ and enhanced 2DEG densities and mobilities were grown by molecular beam epitaxy on InP substrates. The increased conduction band offset resulted in extremely high electron densities of $3.81 \times 10^{12} \text{ cm}^{-2}$, with a 4.2 K mobility of $51100 \text{ cm}^2/\text{Vs}$ for single-sided doping and highest ever densities of $6.70 \times 10^{12} \text{ cm}^{-2}$ for double-sided doping. The strain-symmetrization with stress compensating $\text{In}_y\text{Al}_{1-y}\text{As}$ barriers did not, however, allow an increase in the critical In content of the $\text{In}_x\text{Ga}_{1-x}\text{As}$ QW.

1. Introduction

InGaAs/InAlAs high electron mobility transistors (HEMTs) on InP substrates are currently the most promising devices for high frequency amplification [1-5]. Device performance can be improved either by reducing the gate length [1,3] or by increasing the channel conductance [3,6-9]. The enhanced conduction band offset due to the increasing In content in the active channel leads to higher electron mobilities and saturation velocities. Numerous investigations have been made to maximize the In content in the channel [2,7,8], including stress compensated structures with opposite strain in quantum wells and barriers. We investigated strain-symmetrized SQW structures with a homogeneously and compressively strained well and $\text{In}_y\text{Al}_{1-y}\text{As}$ under tensile strain for stress compensation. The essential advantages of these structures are a larger conduction band discontinuity due to the increased band gap in the $\text{In}_y\text{Al}_{1-y}\text{As}$ barrier, and a better pinch-off of the HEMTs due to improved electron confinement in the 12 nm wide QW. We demonstrate the optimization of 2DEG mobility and density in these structures with modified In contents in the barriers and the wells and, secondly, differing doping profiles. Optimized HEMT structures with single-sided and double-sided doping yielding 2DEG densities, measured at 4.2 K, as high as $3.81 \times 10^{12} \text{ cm}^{-2}$ and $6.70 \times 10^{12} \text{ cm}^{-2}$, respectively.

2. MBE growth

All structures used here were grown by solid source MBE on (100) oriented InP:Fe substrates with growth rates around 1 $\mu\text{m/h}$, a substrate temperature of 530 $^{\circ}\text{C}$ and a V/III ratio (beam equivalent pressure) of 27. The different compositions of the ternary layers were obtained by controlling the group III flux ratio.

Figure 1 schematically shows the epitaxial layer sequence of the investigated HEMT structure. The buffer consists of two 100 nm thick $\text{In}_{0.52}\text{Al}_{0.48}\text{As}$ layers and two $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}/\text{In}_{0.52}\text{Al}_{0.48}\text{As}$ superlattices. The temperatures of the Ga effusion cell was decreased while growing the second $\text{In}_{0.52}\text{Al}_{0.48}\text{As}$ layer from the value necessary for lattice matched growth to the value appropriate for the strained QW. The In flux was kept constant, except for the growth of the lattice matched cap layer. As our MBE machine is equipped with two Al cells (and a single Ga and In cell each), we were able to grow the layer sequence shown in Fig. 1 without any growth interruption at the QW interfaces, having different Al fluxes at our disposal.

The investigated range in the In contents of the QWs and the barriers are given in Table I. The Si doping level in the $\text{In}_y\text{Al}_{1-y}\text{As}$ supply layer was $1.5 \times 10^{19} \text{ cm}^{-3}$ for sample M5409, M5424 and M5425 and $1.0 \times 10^{19} \text{ cm}^{-3}$ for all others. The measured transport properties reflect the influence of the parallel conduction in the cap layers: Samples M5406 - M5410 have doped caps of only 7 nm thickness, whereas the other samples have thicker caps (5 nm undoped InGaAs topped by 10 nm doped InGaAs). The doping in the cap layers is $3.0 \times 10^{18} \text{ cm}^{-3}$ in all samples.

The characterization of our HEMT structures was made by Hall measurements at 300, 77 and 4.2 K and Shubnikov-de Haas (SdH) measurements at 4.2 K to separate the 2DEG conductance from the parallel conductance in the cap layers. The thicknesses of the surface

InGaAs	$3 \cdot 10^{18} \text{ cm}^{-3}$	$y=0.53$	10 nm
InGaAs	-	$y=0.53$	5 nm
InAlAs	-	$y=0.52$	7 nm
InAlAs	-	$y=0.52 - 0.415$	3 nm
InAlAs	$1 - 1.5 \cdot 10^{19} \text{ cm}^{-3}$	$y=0.52 - 0.415$	12.5 nm
InAlAs	-	$y=0.52 - 0.415$	2 nm
InGaAs	-	$y=0.53 - 0.74$	12 nm
InAlAs	-	$y=0.52 - 0.415$	6 nm
InAlAs	$0 - 1.5 \cdot 10^{19} \text{ cm}^{-3}$	$y=0.52 - 0.415$	2 nm
InAlAs	-	$y=0.52 - 0.415$	16 nm
InAlAs	-	$y=0.52$	100 nm
SL			
InAlAs	-	$y=0.52$	100 nm
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InP (:Fe) Substrat			

Figure 1: Schematic layer structure of a strain-symmetrized HEMT.

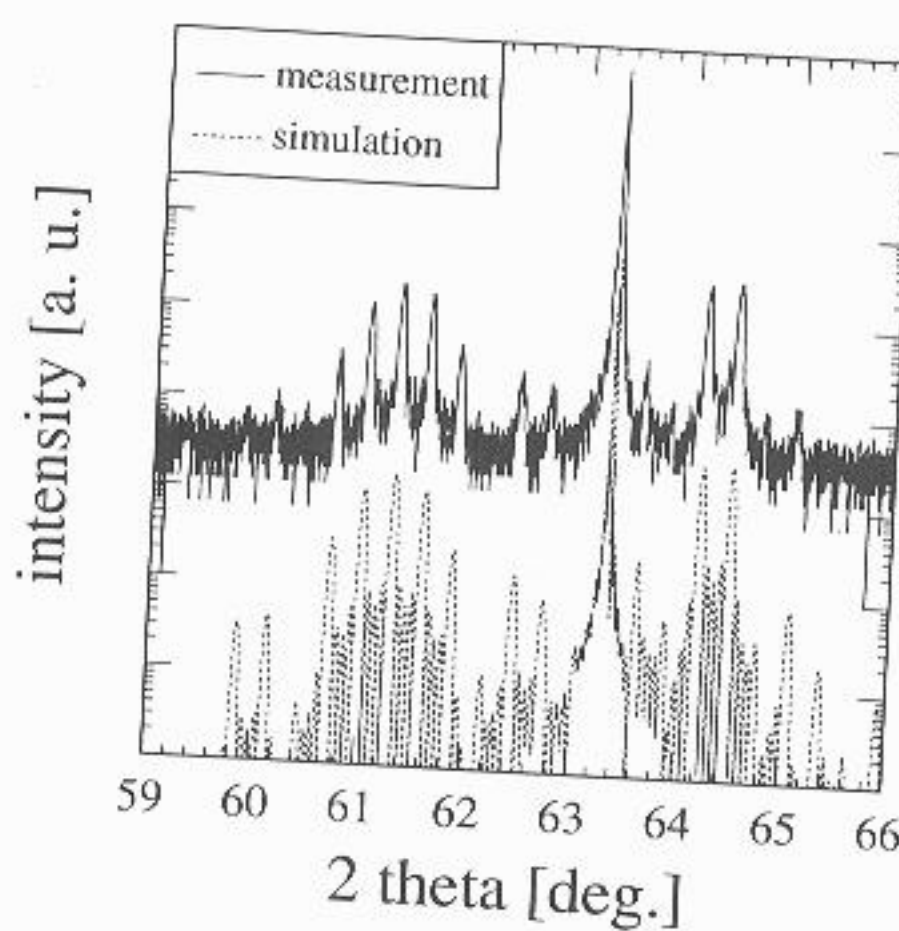


Figure 2: X-ray diffraction spectra (400) of a strain-symmetrized MQW structure (5 periods) with 1 μm $\text{In}_{0.52}\text{Al}_{0.48}\text{As}$ buffer layer on InP substrate (barrier: 23.8 nm $\text{In}_{0.53}\text{Al}_{0.47}\text{As}$, QW: 12.0 nm $\text{In}_{0.749}\text{Ga}_{0.251}\text{As}$; solid: measurement, dotted: dynamical scattering theory).

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3. Structural

Figure 2 sho
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Table I: Structural parameters of HEMT structures

Sample

x ($\text{In}_x\text{Ga}_{1-x}\text{As}$) QW
y ($\text{In}_y\text{Al}_{1-y}\text{As}$) Barr
doping N_D [10^{19} cm^{-3}]
$\mu_{\text{Hall}}(300\text{K})$ [cm^2/Vs]
$n_{\text{Hall}}(300\text{K})$ [10^{12} cm^{-2}]
$\mu_{\text{Hall}}(77\text{K})$ [cm^2/Vs]
$n_{\text{Hall}}(77\text{K})$ [10^{12} cm^{-2}]
$\mu_{\text{Hall}}(4.2\text{K})$ [cm^2/Vs]
$n_{\text{Hall}}(4.2\text{K})$ [10^{12} cm^{-2}]
$n_{1,\text{SdH}}(4.2\text{K})$ [10^{12} cm^{-2}]
$n_{2,\text{SdH}}(4.2\text{K})$ [10^{12} cm^{-2}]
LM=lattice matched
ssd=single-sided dop

on (100) oriented InP:Fe substrates at a temperature of 530 °C and a V/III ratio (beam flux) of 10¹⁹ cm⁻². The compositions of the ternary layers were obtained by

sequence of the investigated HEMT structure. The nominal layer sequence is a 1 µm thick lattice matched In_{0.52}Al_{0.48}As buffer layer, followed by 5 QWs, having 24 nm thick In_{0.415}Al_{0.585}As barriers and 12 nm wide In_{0.74}Ga_{0.26}As wells. The simulated spectrum using dynamical scattering theory fits the measured curve excellently, with only small deviations from the nominal values. The compositions and thicknesses from the simulation are: 23.8 nm In_{0.419}Al_{0.581}As, 12.0 nm In_{0.749}Ga_{0.251}As, 1 µm In_{0.522}Al_{0.478}As buffer. Although we used high growth temperatures of 530 °C and a wide range of In contents from 0.415 to 0.74 in the QW and the barriers, the simulation indicates that there is no noticeable In segregation at the interfaces, in contrast to observed segregation lengths of 3.5 nm at the interfaces of pseudomorphic In_{0.74}Ga_{0.26}As/In_{0.52}Al_{0.48}As MQWs with lattice matched barriers, which were grown under identical growth conditions [13]. In segregation has also been reported for the growth of strained InGaAs layers on GaAs [10-12].

QWs and the barriers are given in Table I. The carrier density in the channel layer was 1.5 × 10¹⁹ cm⁻³ for sample M5406 and M5407, and 1.0 × 10¹⁹ cm⁻³ for all others. The measured transport properties are given in Table II. The measured transport properties in the cap layers: Samples M5406 and M5407 have a carrier density of 1.5 × 10¹⁹ cm⁻³, whereas the other samples have a carrier density of 1.0 × 10¹⁹ cm⁻³ (undoped InGaAs). The doping in the cap layers was controlled by Hall measurements at 300 K, and the carrier density was measured at 4.2 K to separate the 2DEG and the bulk carrier densities. The thicknesses of the surface layers were controlled by X-ray reflectivity measurements (XRR). Compositions of the strained layers were monitored by X-ray diffraction (XRD) on strain-symmetrized MQW structures, always grown in direct sequence with the HEMT structures.

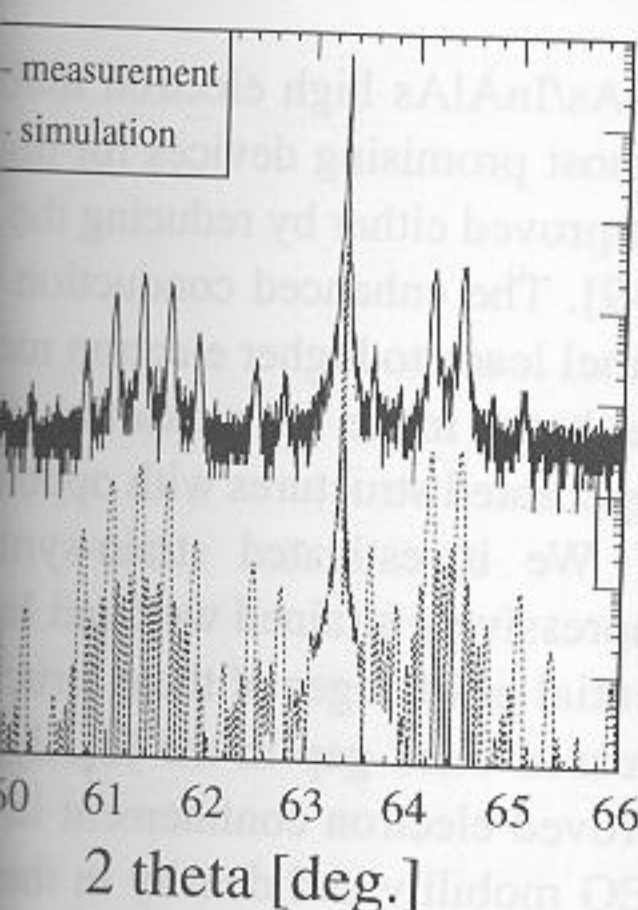
3. Structural and transport properties

Figure 2 shows the XRD spectra of a strain-symmetrized In_{0.74}Ga_{0.26}As/In_{0.415}Al_{0.585}As MQW structure. The nominal layer sequence is a 1 µm thick lattice matched In_{0.52}Al_{0.48}As buffer layer, followed by 5 QWs, having 24 nm thick In_{0.415}Al_{0.585}As barriers and 12 nm wide In_{0.74}Ga_{0.26}As wells. The simulated spectrum using dynamical scattering theory fits the measured curve excellently, with only small deviations from the nominal values. The compositions and thicknesses from the simulation are: 23.8 nm In_{0.419}Al_{0.581}As, 12.0 nm In_{0.749}Ga_{0.251}As, 1 µm In_{0.522}Al_{0.478}As buffer. Although we used high growth temperatures of 530 °C and a wide range of In contents from 0.415 to 0.74 in the QW and the barriers, the simulation indicates that there is no noticeable In segregation at the interfaces, in contrast to observed segregation lengths of 3.5 nm at the interfaces of pseudomorphic In_{0.74}Ga_{0.26}As/In_{0.52}Al_{0.48}As MQWs with lattice matched barriers, which were grown under identical growth conditions [13]. In segregation has also been reported for the growth of strained InGaAs layers on GaAs [10-12].

Table I: Structural parameters, mobilities and 2-DEG densities of strain-symmetrized In_xGa_{1-x}As/In_yAl_{1-y}As HEMT structures with different In-content in the channel and the barriers (QW thickness: 12 nm)

Sample	LM	PM	PM	SS	SS	PM	SS	PM	SS
	ssd	ssd	ssd	ssd	ssd	ssd	ssd	dsd	dsd
	M5406	M5407	M5408	M5420	M5421	M5409	M5424	M5410	M5425
x (In _x Ga _{1-x} As) QW	0.53	0.67	0.74	0.67	0.74	0.74	0.74	0.74	0.74
y (In _y Al _{1-y} As) Barrier	0.52	0.52	0.52	0.45	0.415	0.52	0.415	0.52	0.415
doping N _D [10 ¹⁹ cm ⁻³]	1.0	1.0	1.0	1.0	1.0	1.5	1.5	1.0	1.5
μ _{Hall} (300K) [cm ² /Vs]	9420	11390	12220	11040	12150	11090	11150	7450	7640
n _{Hall} (300K) [10 ¹² cm ⁻²]	3.16	3.36	3.49	3.72	3.82	4.22	4.56	6.58	7.67
μ _{Hall} (77K) [cm ² /Vs]	30960	43000	50220	43150	50070	40700	46860	14160	16280
n _{Hall} (77K) [10 ¹² cm ⁻²]	3.03	3.21	3.35	3.53	3.67	4.05	4.25	6.55	7.31
μ _{Hall} (4.2K) [cm ² /Vs]	33540	54940	62930	46820	-	54970	51080	16230	15060
n _{Hall} (4.2K) [10 ¹² cm ⁻²]	2.79	2.97	3.18	3.79	-	3.64	4.30	6.07	6.92
n _{1,SdH} (4.2K) [10 ¹² cm ⁻²]	2.55	2.68	2.74	2.81	-	3.05	3.12	3.84	4.15
n _{2,SdH} (4.2K) [10 ¹² cm ⁻²]	0.26	0.38	0.45	0.46	-	0.61	0.69	2.20	2.55

LM=lattice matched QW and barriers; PM=pseudomorphic QW; SS=strain-symmetrized structure
ssd=single-sided doping in top barrier; dsd= double-sided doping in both barriers



X-ray diffraction spectra (400) of a strain-symmetrized MQW structure (5 periods) with 12 nm In_{0.53}Al_{0.47}As, QW: 12.0 nm In_{0.74}Ga_{0.26}As; solid: measurement, dotted: dynamical scattering theory).

The transport properties of all the HEMT structures investigated in the course of this work are given in Table I. Sample M5406 is a lattice matched reference sample with the QW width of 12 nm used in all our HEMT structures. The mobilities of this reference structure were measured to $9400 \text{ cm}^2/\text{Vs}$ and $33500 \text{ cm}^2/\text{Vs}$ at 300 and 4.2 K, respectively. These mobilities are clearly lower than those obtained on lattice matched structures with 32 nm wide QWs ($10000 \text{ cm}^2/\text{Vs}$ and $42000 \text{ cm}^2/\text{Vs}$) due to increased interface roughness scattering by the bottom barrier.

Using compressive strain only in the QWs, i.e. with lattice matched barriers, increases the mobilities with rising In content from $11400 \text{ cm}^2/\text{Vs}$ and $55000 \text{ cm}^2/\text{Vs}$ for $x = 0.67$ (sample M5407) to $12200 \text{ cm}^2/\text{Vs}$ and $63000 \text{ cm}^2/\text{Vs}$ for $x = 0.74$ (sample M5408). Simultaneously there is an increase in the 4.2 K 2DEG densities from $2.81 \times 10^{12} \text{ cm}^{-2}$ to $3.06 \times 10^{12} \text{ cm}^{-2}$ and $3.19 \times 10^{12} \text{ cm}^{-2}$ (samples M5406, M5407 and M5408). SdH measurements clearly show two occupied subbands with densities n_1 and n_2 . The intersubband separation in sample M5406 was found to be 108 meV by the peak separation in photoluminescence, agreeing well with the observed values of n_1 and n_2 using the calculated density of states.

On increasing the doping level to $1.5 \times 10^{19} \text{ cm}^{-3}$ in the supply layer, the total 2DEG density reaches $3.66 \times 10^{12} \text{ cm}^{-2}$, while still maintaining a 4.2 K mobility of $55000 \text{ cm}^2/\text{Vs}$ (sample M5409, lattice matched barriers, single-sided doping). Using $\text{In}_y\text{Al}_{1-y}\text{As}$ barriers under tensile strain increases the carrier density even further. Sample M5424 is strain-symmetrized with $x = 0.74$ in the QW and $y = 0.415$ in the barriers, showing a total 2DEG density of $3.81 \times 10^{12} \text{ cm}^{-2}$ at 4.2 K. The mobilities are practically unchanged (see Table I).

A further increase of the In content in the QW to $x = 0.78$ with a corresponding decrease in the barrier to $x = 0.34$, however, resulted in a drastic decrease in mobility due to

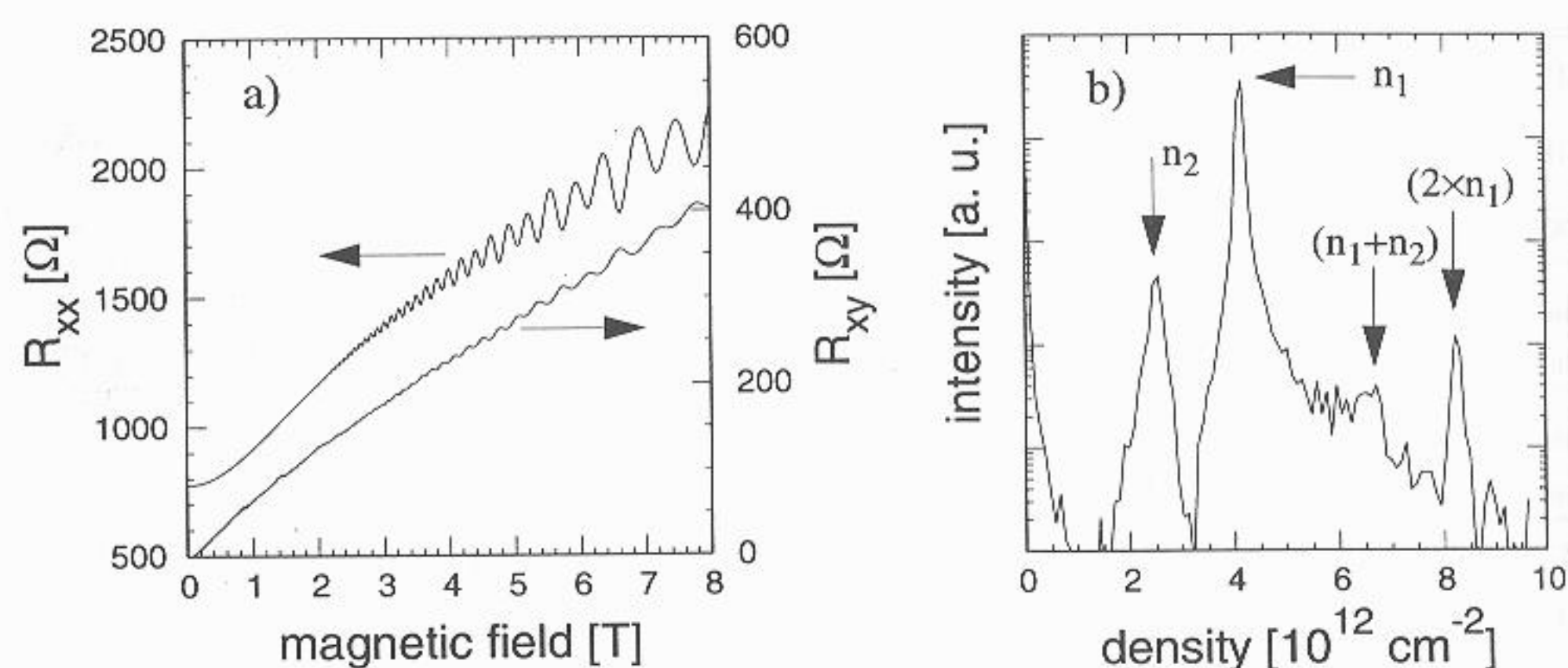


Figure 3: a) Shubnikov-de Haas measurement on a strain-symmetrized HEMT structure with double-sided doping (sample M5425); b) Fourier transformation of R_{xx} .

relaxation and the formation of critical In content of strain compensating layers, on content is the same in ps

Extremely high 2DEG densities from both - top and bottom barriers of $6.70 \times 10^{12} \text{ cm}^{-2}$ were found respectively. The first spacer of 4 nm, the conduction 2DEG densities are, to compare. The marked reduction in single-sided doping is observed in structure M5425 we found respectively. In spite of this we were able to determine measurements and Fourier of 4.15 and 2.55×10^{12}

4. High electron mobility

HEMT devices with a gate using e-beam lithography increasing In content from were lattice matched transconductance from intrinsic values from 10 gain were measured to

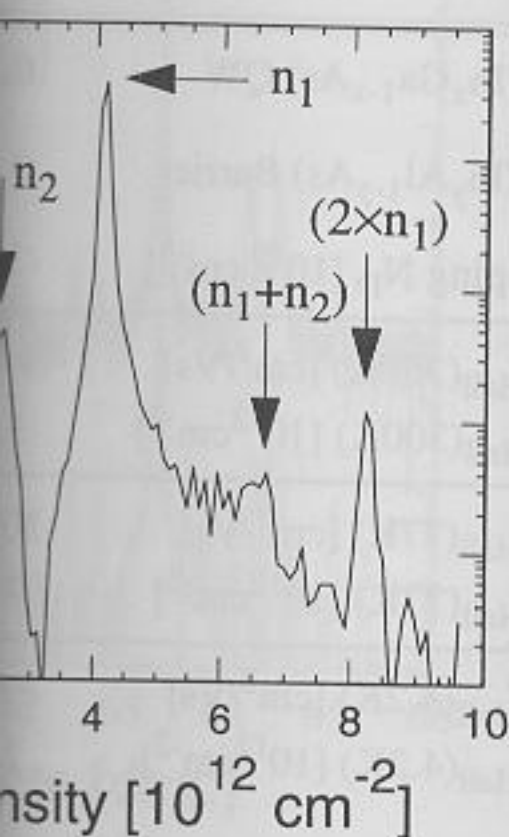
5. Summary

Single-sided doped structures showed 2DEG densities. Extremely high 2DEG densities from both barriers. XRD of In at the interfaces was

6. Acknowledgement

We thank the German Research Foundation (DFG) for financial support

in the course of this work are sample with the QW width of is reference structure were respectively. These mobilities res with 32 nm wide QWs roughness scattering by the matched barriers, increases $55000 \text{ cm}^2/\text{Vs}$ for $x = 0.67$ $x = 0.74$ (sample M5408). s from $2.81 \times 10^{12} \text{ cm}^{-2}$ to 5407 and M5408). SdH densities n_1 and n_2 . The neV by the peak separation s of n_1 and n_2 using the pply layer, the total 2DEG mobility of $55000 \text{ cm}^2/\text{Vs}$ Using $\text{In}_y\text{Al}_{1-y}\text{As}$ barriers ther. Sample M5424 is e barriers, showing a total ractically unchanged (see 0.78 with a corresponding decrease in mobility due to



structure with double-sided

relaxation and the formation of misfit dislocations. Whereas it has been reported [8] that the critical In content of strained InGaAs QWs could be increased by the insertion of stress compensating layers, our results show that, under our growth conditions, the maximum In content is the same in pseudomorphic and strain-symmetrized structures.

Extremely high 2DEG densities were, however, obtained using double-sided doping from both - top and bottom - barriers. Total 2DEG densities as high as $6.04 \times 10^{12} \text{ cm}^{-2}$ and $6.70 \times 10^{12} \text{ cm}^{-2}$ were found by SdH measurements at 4.2 K in samples M5410 and M5425, respectively. The first sample is doped to $1.0 \times 10^{19} \text{ cm}^{-3}$ in the supply layer with a bottom spacer of 4 nm, the corresponding values for the latter are $1.5 \times 10^{19} \text{ cm}^{-3}$ and 6 nm. These 2DEG densities are, to our knowledge, the highest ever reported for InGaAs/InAlAs HEMTs. The marked reduction in mobility of double-sided doped structures in comparison to single-sided doping is due to Si segregation towards the channel [14]. So, for example in structure M5425 we found mobilities of $7600 \text{ cm}^2/\text{Vs}$ and $15100 \text{ cm}^2/\text{Vs}$ at 300 and 4.2 K, respectively. In spite of the low mobility and considerable parallel conduction in the cap layer we were able to determine clearly the separate densities of the two occupied subbands by SdH measurements and Fourier transformation, as shown in Fig. 3. This analysis yields densities of 4.15 and $2.55 \times 10^{12} \text{ cm}^{-2}$ for the first and the second subband in the QW.

4. High electron mobility transistors

HEMT devices with a gate length of $0.13 \mu\text{m}$ were fabricated from sample M5406 - M5408, using e-beam lithography and wet chemical selective gate recess with succinic acid. The increasing In content from $x = 0.53$ to 0.74 in the active QW (the barriers of these samples were lattice matched to the InP substrate) resulted in an increase in the extrinsic transconductance from 670 mS/mm to 1150 mS/mm , with a concurrent increase in the intrinsic values from 1000 mS/mm to 1600 mS/mm . The cut-off frequencies for the current gain were measured to $200 - 220 \text{ GHz}$ for all devices.

5. Summary

Single-sided doped strain-symmetrized $\text{In}_{0.74}\text{Ga}_{0.26}\text{As}/\text{In}_{0.415}\text{Al}_{0.585}\text{As}$ HEMT structures showed 2DEG densities as high as $3.81 \times 10^{12} \text{ cm}^{-2}$ and mobilities of $51100 \text{ cm}^2/\text{Vs}$ at 4.2 K. Extremely high 2DEG densities of $6.70 \times 10^{12} \text{ cm}^{-2}$ were obtained using double-sided doping from both barriers. XRD of strain-symmetrized MQW structures showed that the segregation of In at the interfaces was significantly suppressed by strain-symmetrization.

6. Acknowledgement

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